

Electromagnetic radiation.

Homework

1. It can be shown that the following electric and magnetic fields,

$$\begin{aligned} E_r &= ke^{i(\omega t - kr)} \left[-\frac{1}{(kr)^2} + \frac{i}{(kr)^3} \right] \cos \theta, \\ E_\theta &= \frac{1}{2} ke^{i(\omega t - kr)} \left[-\frac{i}{kr} - \frac{1}{(kr)^2} + \frac{i}{(kr)^3} \right] \sin \theta, \\ H_\psi &= \frac{1}{2} k \sqrt{\frac{\epsilon_0}{\mu_0}} e^{i(\omega t - kr)} \left[-\frac{i}{kr} - \frac{1}{(kr)^2} \right] \sin \theta, \\ E_\psi &= H_r = H_\theta = 0, \end{aligned}$$

written in spherical polar coordinates (r, θ, ψ) , with $\omega/k = 1/\sqrt{\epsilon_0\mu_0}$, satisfy Maxwell's equations in free space in the absence of charges and currents.

- (a) Consider an oscillating electric dipole of moment $\mathbf{P}e^{i\omega t}$, taking \mathbf{P} along the z -axis. Show by considering the near-field ($kr \ll 1$) limit that the fields due to the oscillating electric dipole can be obtained by multiplying the above solution of Maxwell's equations by

$$-\frac{ik^2\mathbf{P}}{2\pi\epsilon_0}.$$

- (b) Obtain an expression valid in the far-field ($kr \gg 1$) limit for the time-averaged Poynting vector due to such an oscillating electric dipole at the origin and show that the total flux of radiation outward through a sphere is given by

$$\frac{\mu_0\sqrt{\epsilon_0\mu_0}\omega^4P^2}{12\pi}.$$

[Exam 2003 Question 7]

2. (a) By making a suitable expansion of the differentials $d[(\mathbf{b} \cdot \mathbf{r})\mathbf{r}]$ and $\mathbf{b} \times (\mathbf{r} \times d\mathbf{r})$, or otherwise, prove that

$$\oint_C (\mathbf{b} \cdot \mathbf{r}) d\mathbf{r} = \mathbf{M} \times \mathbf{b},$$

where

$$\mathbf{M} \equiv \frac{1}{2} \oint_C \mathbf{r} \times d\mathbf{r},$$

C is a closed curve, and \mathbf{b} is a constant vector. If C is a plane curve, show that the magnitude of \mathbf{M} is equal to the area enclosed by C , and that the direction of \mathbf{M} is perpendicular to the plane of C , and related to the right-hand screw to the direction in which the curve is traversed.

- (b) A circular loop of wire of radius a carries an alternating current $I = I_0 \cos \omega t$, where I_0 and ω are constants and $\omega a \ll c$, c being the velocity of light. The loop lies in the x - y plane and is centred upon the origin. Show that, correct to first-order in $1/r$, where r is the distance from the origin, the (complex) vector potential is given by

$$\mathbf{A}(\mathbf{r}, t) = \frac{i\mu_0 I_0 a^2 k}{4} \frac{e^{i(\omega t - kr)}}{r} \sin \theta \hat{\psi},$$

where $k \equiv \omega/c$ and (r, θ, ψ) are the polar coordinates of \mathbf{r} . Hence, determine the magnetic field \mathbf{H} and the electric field \mathbf{E} , correct to first order in $1/r$, and the corresponding Poynting vector \mathbf{S} .

3. (a) Show, from Maxwell's equations, how the \mathbf{B} and \mathbf{E} fields may be parameterised in terms of a vector potential \mathbf{A} and a scalar potential ϕ .
- (b) In Lorenz gauge (i.e., $\nabla \cdot \mathbf{A} + \epsilon_0 \mu_0 \partial \phi / \partial t = 0$), the vector potential \mathbf{A} at a large distance r from a finite system of charges is approximately

$$\mathbf{A} = \frac{\mu_0}{4\pi r} \dot{\mathbf{P}}(t - r/c),$$

where $\mathbf{P}(t)$ is the dipole moment of the system, $c = 1/\sqrt{\epsilon_0 \mu_0}$ and $\dot{\mathbf{P}}(t) \equiv d\mathbf{P}/dt$.

For

$$\mathbf{P}(t) = \mathbf{P}_0 e^{i\omega t},$$

where \mathbf{P}_0 and ω are constants, show that, correct to first order in $1/r$, the corresponding magnetic and electric fields are

$$\mathbf{B} = \frac{\omega \mu_0 k}{4\pi} \frac{e^{i(\omega t - kr)}}{r} \hat{\mathbf{r}} \times \mathbf{P}_0,$$

$$\mathbf{E} = \frac{\omega^2 \mu_0}{4\pi} \frac{e^{i(\omega t - kr)}}{r} \hat{\mathbf{r}} \times (\mathbf{P}_0 \times \hat{\mathbf{r}}),$$

where $k = \omega/c$ and $\hat{\mathbf{r}}$ is a unit vector in the direction of \mathbf{r} .

Hint: in deriving the result, the following vector identities may be of use:

$$\nabla \times (f\mathbf{a}) = f(\nabla \times \mathbf{a}) + \nabla f \times \mathbf{a},$$

$$\nabla \cdot (f\mathbf{a}) = f(\nabla \cdot \mathbf{a}) + \nabla f \cdot \mathbf{a},$$

$$\nabla(fg) = f\nabla g + g\nabla f.$$

- (c) A rigid electric dipole p is centred on the origin and lies in the x - y plane. The dipole rotates anticlockwise with constant angular velocity ω , and is along the x axis at $t = 0$. Show that, correct to first order in $1/r$, the \mathbf{B} and \mathbf{E} fields created by the dipole are:

$$\mathbf{B} = \frac{p\omega\mu_0 k}{4\pi r} [\cos(\omega t - kr - \psi) \cos \theta \hat{\boldsymbol{\psi}} - \sin(\omega t - kr - \psi) \hat{\boldsymbol{\theta}}],$$

$$\mathbf{E} = \frac{pk^2}{4\pi\epsilon_0 r} [\sin(\omega t - kr - \psi) \hat{\boldsymbol{\psi}} + \cos(\omega t - kr - \psi) \cos \theta \hat{\boldsymbol{\theta}}],$$

where (r, θ, ψ) are spherical polar coordinates.

Find the corresponding Poynting vector.

Further Examples

1. (a) A set of point charges is in motion in a finite region surrounding the origin. If the speeds of the charges are much smaller than the velocity of light c , show that at large distance r from the origin the vector potential \mathbf{A} is given by

$$\mathbf{A} \simeq \frac{\mu_0}{4\pi r} \dot{\mathbf{P}}(t - r/c),$$

where $\mathbf{P}(t)$ is the dipole moment of the system at time t .

- (b) For $\mathbf{P}(t) = \mathbf{P}_0 e^{i\omega t}$, where \mathbf{P}_0 and ω are constants, find, correct to first order in $1/r$, the corresponding electric and magnetic fields at large r .

$$\mathbf{E} = \frac{\omega^2 \mu_0}{4\pi} \frac{e^{i(\omega t - kr)}}{r} \hat{\mathbf{r}} \times (\mathbf{P}_0 \times \hat{\mathbf{r}}), \quad \mathbf{B} = \frac{\omega \mu_0 k}{4\pi} \frac{e^{i(\omega t - kr)}}{r} \hat{\mathbf{r}} \times \mathbf{P}_0.$$

Answers