

Examples

①

$$\frac{d^2y}{dx^2} + \lambda y = 0, \quad y'(0) = y(l) = 0.$$

This is a 2nd-order linear differential equation with constant coefficients. To solve it, seek solution in the form

$$y = e^{sx}$$

$$y' = se^{sx}, \quad y'' = s^2 e^{sx}$$

Substitute into the equation:

$$s^2 e^{sx} + \lambda e^{sx} = 0 \quad \left\{ \text{multiply by } e^{-sx} \right.$$

$$s^2 + \lambda = 0 : \underline{\text{auxiliary equation}}$$

$$s^2 = -\lambda$$

$$s = \pm \sqrt{-\lambda}$$

If $\lambda \neq 0$, the auxiliary equation has two distinct roots, and the general solution of the DE is:

$$y = C_1 e^{\sqrt{-\lambda}x} + C_2 e^{-\sqrt{-\lambda}x},$$

where C_1 and C_2 are arbitrary constants.

To implement the boundary conditions,

$$y' = \sqrt{-\lambda} C_1 e^{\sqrt{-\lambda}x} - C_2 \sqrt{-\lambda} e^{-\sqrt{-\lambda}x}$$

$$y'(0) = 0 \text{ gives: } \sqrt{-\lambda} (C_1 - C_2) = 0$$

$$\Rightarrow C_2 = C_1 \equiv C. \quad (\lambda \neq 0).$$

$$\text{So } y = C(e^{\sqrt{-\lambda}x} + e^{-\sqrt{-\lambda}x}) \quad (2)$$

Using the 2nd boundary condition $y(l)=0$:

$$C(e^{\sqrt{-\lambda}l} + e^{-\sqrt{-\lambda}l}) = 0$$

For $\lambda < 0$ the sum of two real exponents cannot be zero $\Rightarrow C=0$ and the equation has only a trivial solution $y=0$.

Consider $\lambda > 0$. Then $\sqrt{-\lambda} = i\sqrt{\lambda}$

and $y = C(e^{i\sqrt{\lambda}x} + e^{-i\sqrt{\lambda}x})$

Recall:

$e^{i\theta} = \cos\theta + i\sin\theta$

$e^{-i\theta} = \cos\theta - i\sin\theta$

$\cos\theta = \frac{e^{i\theta} + e^{-i\theta}}{2}$

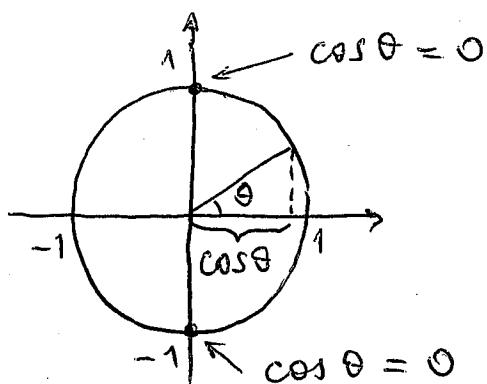
$\sin\theta = \frac{e^{i\theta} - e^{-i\theta}}{2i}$

$y = (2C) \cos(\sqrt{\lambda}x)$
rename "C"

$y = C \cos(\sqrt{\lambda}x)$

Now, the boundary condition

requires: $C \cos(\sqrt{\lambda}l) = 0 \quad (*)$



$$\cos\theta = 0 \quad \text{for}$$

$$\theta = \frac{\pi}{2} + \pi n, \quad n \in \mathbb{Z}.$$

Hence, (*) gives:

$$\sqrt{\lambda}l = \frac{\pi}{2} + \pi n, \quad n=0, 1, \dots$$

$$\sqrt{\lambda} = \frac{\pi(n+\frac{1}{2})}{l} \quad (\sqrt{\lambda}l > 0).$$

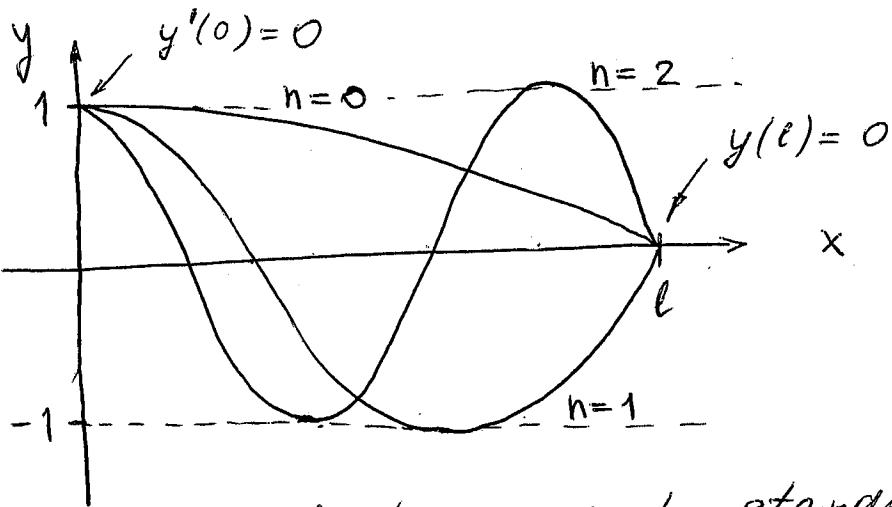
$$\boxed{\lambda = \frac{\pi^2(n+\frac{1}{2})^2}{l^2}}$$

Values of λ for which the equation has nontrivial solutions.

The corresponding solutions are : (3)

$$y = C \cos \left[\pi \left(n + \frac{1}{2} \right) \frac{x}{l} \right]$$

We can sketch the first few (for $C=1$)



These solutions in fact represent standing waves, e.g., those that occur when air vibrates in a semi-closed (organ) pipe:

$y(x)$ then represents the displacement of air along the pipe.

Naturally $y(l) = 0$
(at the closed end)

and y is a maximum
($y'(0) = 0$) near the open end,

$n=0$

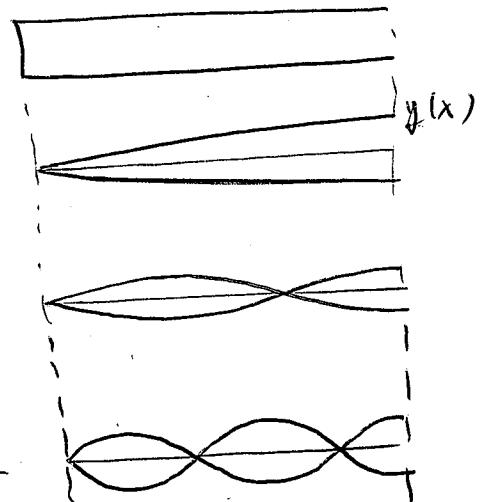
$n=1$

$n=2$

etc.

$x=l$

$x=0$



$$② u = f(x^2 - y^2) \quad (*)$$

By taking derivatives of u , with respect to x and y , we should be able to eliminate f and obtain a partial DE, for which $(*)$ is the solution.

$$(1) u_x = \frac{\partial u}{\partial x} = f'(x^2 - y^2) \cdot 2x \quad \left. \begin{array}{l} \text{using chain} \\ \text{rule} \end{array} \right\}$$

[f' is the derivative of f with respect to its argument]

$$(2) u_y = \frac{\partial u}{\partial y} = f'(x^2 - y^2) (-2y)$$

Dividing (1) by (2), we have:

$$\frac{u_x}{u_y} = - \frac{2x}{2y}$$

$$yu_x = -xu_y$$

$$\underline{yu_x + xu_y = 0}$$

$$③ u_{xxx} + u_x = 0 \quad \text{for } u(x, y, z)$$

Integrating this equation with respect to x , we have:

$$u_{xx} + u = f(y, z) \quad (*)$$

[The right-hand side is an arbitrary "constant" (with respect to x), i.e. an arbitrary function of y and z .]

$(*)$ is a linear inhomogeneous equation.

Solving the homogeneous version first, (5)

$$u_{xx} + u = 0,$$

we have:

$$u = g(y, z) \cos x + h(y, z) \sin x$$

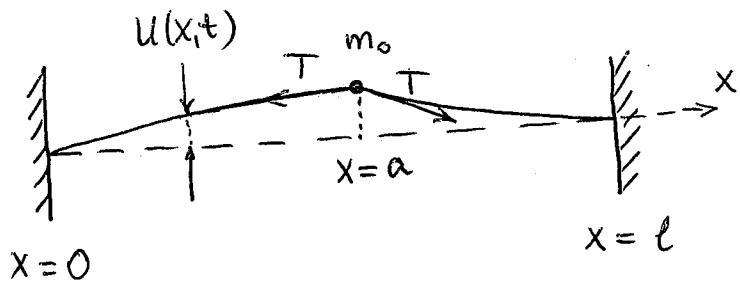
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arbitrary functions of y and z
(i.e. "constants" with respect to x).

To solve the inhomogeneous equation (*), it sufficient to add $f(y, z)$ to the above solution.

$$\underline{u(x, y, z) = f(y, z) + g(y, z) \cos x + h(y, z) \sin x}$$

(one can check that this is correct by back-substitution into (*).)

(4)



Boundary conditions:

$$u(0, t) = 0$$

$$u(l, t) = 0$$

$u(x, t)$ - displacement of the string.

For $0 \leq x < a$ and $a < x \leq l$ the wave equation is exactly the same as that for a uniform string,

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0 \quad (*)$$

$$c = \sqrt{\frac{T}{\rho}}$$

T - tension

ρ - linear mass density

Force acting on m_0 :

$$T \sin \theta|_{x=a+0} - T \sin \theta|_{x=a-0}$$

Here $x = a + 0$ and $x = a - 0$ mean values
of x taken "just to the right" and "just to the left"
of a . The angle θ is that the string makes with
the x axis. For small θ , $\theta \ll 1$,

$$\sin \theta \approx \tan \theta = \frac{\partial u}{\partial x}.$$

Hence, the force acting on m_0 is $T \left(\frac{\partial u}{\partial x} \Big|_{x=a+0} - \frac{\partial u}{\partial x} \Big|_{x=a-0} \right)$

Newton's 2nd law for m_0 :

$$m_0 \frac{\partial^2 u}{\partial t^2} \Big|_{x=a} = T \left(\frac{\partial u}{\partial x} \Big|_{x=a+0} - \frac{\partial u}{\partial x} \Big|_{x=a-0} \right) \quad (*)$$

Therefore, in order to solve the problem,
one needs to find a solution of $(*)$ for $0 \leq x \leq a$
satisfying $u(0, t) = 0$; let's call it $u_1(x, t)$.
Then do the same for $a \leq x \leq l$ with $u(l, t) = 0$;
call this solution $u_2(x, t)$. The two solutions
must match at $x = a$:

$$u_1(a, t) = u_2(a, t),$$

and $(**)$ must be satisfied:

$$\frac{m_0}{T} \frac{\partial^2 u}{\partial t^2} \Big|_{x=a} = \frac{\partial u_2}{\partial x} \Big|_{x=a} - \frac{\partial u_1}{\partial x} \Big|_{x=a}.$$