

Homework problemsSOLUTIONS

$$\begin{aligned} \textcircled{1} \quad (\text{a}) \quad \mathcal{L}[e^{\alpha t} f(t)] &= \int_0^\infty e^{-pt} e^{\alpha t} f(t) dt \\ &= \int_0^\infty e^{-(p-\alpha)t} f(t) dt \\ &= \underline{F(p-\alpha)}, \end{aligned}$$

where $F(p) = \int_0^\infty e^{-pt} f(t) dt$.

$$\begin{aligned} \text{(b)} \quad \mathcal{L}[f(t-a)] &= \int_0^\infty e^{-pt} f(t-a) dt \\ &= \int_0^\infty e^{-p(t-a)} e^{-pa} f(t-a) dt \\ &= e^{-pa} \int_0^\infty e^{-p(t-a)} f(t-a) dt \end{aligned}$$

$f(t-a) = 0$ for $t-a < 0$, i.e. $t < a$.

Hence, the lower limit can be replaced by a .

Introducing the new variable, $\tau = t-a$, we have:

$$\begin{aligned} \mathcal{L}[f(t-a)] &= e^{-pa} \int_a^\infty e^{-p(t-a)} f(t-a) d(t-a) \\ &= e^{-pa} \int_0^\infty e^{-p\tau} f(\tau) d\tau \\ &= \underline{e^{-pa} F(p)}. \end{aligned}$$

$$\begin{aligned}
 ② \quad \mathcal{L}[y'] &= \int_0^\infty e^{-pt} y'(t) dt \\
 &= \underbrace{e^{-pt} y(t) \Big|_0^\infty}_{\text{Here, the upper limit contribution is zero}} - \int_0^\infty (-p) e^{-pt} y(t) dt \\
 &= -y(0) + p \underbrace{\int_0^\infty e^{-pt} y(t) dt}_{Y(p)}
 \end{aligned}$$

Hence, $\underline{\mathcal{L}[y'] = p Y(p) - y(0)}$.

$$③ \quad y'' - 3y' + 2y = 0, \quad y(0) = 1, \quad y'(0) = -1$$

$$\mathcal{L}[y''] = p^2 Y(p) - p y(0) - y'(0)$$

$$\mathcal{L}[y'] = p Y(p) - y(0)$$

Taking the Laplace transform of the differential equation, and making use of the initial conditions, we have:

$$p^2 Y(p) - p \cdot 1 - (-1) - 3(p Y(p) - 1) + 2 Y(p) = 0$$

$$(p^2 - 3p + 2) Y(p) = p - 4$$

Factorising the left-hand side, we obtain:

$$Y(p) = \frac{p - 4}{(p-1)(p-2)}$$

Using partial fractions, we can expand:

$$\frac{p - 4}{(p-1)(p-2)} = \frac{A}{p-2} + \frac{B}{p-1}$$

$$p-4 = A(p-1) + B(p-2) \quad (3)$$

$$p-4 = (A+B)p - A - 2B$$

$$\text{Hence, } A+B = 1 \Rightarrow B = 1-A$$

$$-4 = -A - 2B$$

$$A + 2(1-A) = 4$$

$$-A = 2 \Rightarrow A = -2$$

$$B = 3$$

Therefore:

$$Y(p) = -\frac{2}{p-2} + \frac{3}{p-1}$$

$$y(t) = \mathcal{L}^{-1}[Y(p)] = \underline{-2e^{2t} + 3e^t}.$$

Here we use
 $\mathcal{L}^{-1}\left[\frac{1}{p-\alpha}\right] = e^{\alpha t}$

$$(4) \quad y'' + 2y' = e^{-t}, \quad y(0) = y'(0) = 0.$$

Applying the Laplace transform and taking into account the initial conditions, we have:

$$p^2 Y(p) + 2p Y(p) = \frac{1}{p+1} \quad \left. \begin{array}{l} \\ \end{array} \right\} \quad \mathcal{L}[e^{\alpha t}] = \frac{1}{p-\alpha}$$

$$p(p+2) Y(p) = \frac{1}{p+1}$$

$$Y(p) = \frac{1}{p(p+1)(p+2)}$$

Let us expand the right-hand side using partial fractions:

$$\frac{1}{p(p+1)(p+2)} = \frac{A}{p} + \frac{B}{p+1} + \frac{C}{p+2}$$

(4)

$$1 = A(p+1)(p+2) + Bp(p+2) + Cp(p+1)$$

$$1 = p^2(A+B+C) + p(3A+2B+C) + 2A$$

$$\begin{cases} A+B+C=0 & (1) \\ 3A+2B+C=0 & (2) \\ 2A=1 & (3) \end{cases}$$

$$(3) \Rightarrow \underline{A = \frac{1}{2}} ; \quad (2)-(1) \text{ gives: } 2A+B=0$$

$$B = -2A \Rightarrow \underline{B = -1}$$

$$\underline{C = -A - B = -\frac{1}{2} + 1 = \frac{1}{2}} .$$

Hence, $\underline{Y(p) = \frac{1}{2p} - \frac{1}{p+1} + \frac{1}{2(p+2)}}$

$$\underline{y(t) = \mathcal{L}^{-1}[Y(p)] = \frac{1}{2} - e^{-t} + \frac{1}{2} e^{-2t}} . \quad \left. \begin{array}{l} \mathcal{L}^{-1}\left[\frac{1}{p}\right] = 1 \\ \text{for } t \geq 0 \\ (\text{and } 0 \text{ for } t < 0) \end{array} \right\}$$

(5) (a) $\underline{y'' + \omega^2 y = f(t)}$

Taking the Laplace transform, we have:

$$\underline{p^2 Y(p) - py(0) - y'(0) + \omega^2 Y(p) = F(p)} ,$$

where $F(p) = \mathcal{L}[f]$.

$$(p^2 + \omega^2) Y(p) = F(p) + y'(0) + py(0)$$

$$\underline{Y(p) = \frac{F(p) + y'(0) + py(0)}{p^2 + \omega^2}} .$$

(5)

$$(6) \text{ For } y(0) = y'(0) = 0$$

$$Y(p) = \frac{F(p)}{p^2 + \omega^2}$$

This can be written as

$$Y(p) = \underbrace{\frac{1}{\omega} F(p)}_{\mathcal{L}[f]} \underbrace{\frac{\omega}{p^2 + \omega^2}}_{\mathcal{L}[\sin \omega t]}$$

The right-hand side is a product of two Laplace transforms.

Using the convolution theorem, we find:

$$y(t) = \frac{1}{\omega} \mathcal{L}^{-1} \left[F(p) \frac{\omega}{p^2 + \omega^2} \right]$$

$$y(t) = \frac{1}{\omega} \int_0^t f(\tau) \sin \omega(t-\tau) d\tau .$$

$$(6) \quad \begin{aligned} y' &= 4y - 2z & y(0) &= 2, \\ z' &= 5y + 2z & z(0) &= -2. \end{aligned}$$

Taking the Laplace transforms of the differential equations, and making use of

$$\mathcal{L}[y'] = p Y(p) - y(0)$$

$$\mathcal{L}[z'] = p Z(p) - z(0) ,$$

we have:

$$\left\{ \begin{array}{l} p Y(p) - 2 = 4 Y(p) - 2 Z(p) \end{array} \right.$$

$$\left. \begin{array}{l} p Z(p) + 2 = 5 Y(p) + 2 Z(p) \end{array} \right.$$

$$\left\{ \begin{array}{l} (p-4) Y(p) + 2 Z(p) = 2 \end{array} \right. \quad (1)$$

$$\left. \begin{array}{l} -5 Y(p) + (p-2) Z(p) = -2 \end{array} \right. \quad (2)$$

Multiplying (1) by $(p-2)$ and (2) by 2 ,
 and subtracting (2) from (1) , we have: (6)

$$(p-4)(p-2) Y(p) + 10 Y(p) = 2(p-2) + 4$$

$$(p^2 - 6p + 8 + 10) Y(p) = 2p$$

$$Y(p) = \frac{2p}{p^2 - 6p + 18} .$$

From (1) :

$$\begin{aligned} Z(p) &= 1 - \frac{1}{2} (p-4) Y(p) \\ &= 1 - \frac{1}{2} (p-4) \frac{2p}{p^2 - 6p + 18} \\ &= \frac{p^2 - 6p + 18 - (p-4)p}{p^2 - 6p + 18} \\ \Rightarrow Z(p) &= \frac{-2p + 18}{p^2 - 6p + 18} . \end{aligned}$$

Re-writing $Y(p)$ and $Z(p)$:

$$Y(p) = \frac{2p}{p^2 - 6p + 9 + 9} = \frac{2p}{(p-3)^2 + 9} = \frac{2(p-3) + 6}{(p-3)^2 + 9}$$

$$Z(p) = \frac{-2p + 18}{p^2 - 6p + 18} = \frac{-2(p-3) - 6 + 18}{(p-3)^2 + 9} = \frac{-2(p-3) + 12}{(p-3)^2 + 9}$$

$$\Rightarrow Y(p) = \frac{2(p-3)}{(p-3)^2 + 9} + \frac{6}{(p-3)^2 + 9} ,$$

$$Z(p) = -\frac{2(p-3)}{(p-3)^2 + 9} + \frac{12}{(p-3)^2 + 9} .$$

Using $\mathcal{L}^{-1} \left[\frac{\omega}{p^2 + \omega^2} \right] = \sin \omega t$ (7)

$$\mathcal{L}^{-1} \left[\frac{p}{p^2 + \omega^2} \right] = \cos \omega t$$

together with the 1st shift theorem,

$$\mathcal{L}^{-1} [F(p-\alpha)] = e^{\alpha t} f(t),$$

we obtain:

$$\underline{y(t)} = \mathcal{L}^{-1}[Y(p)] = \frac{2e^{3t} \cos 3t + 2e^{3t} \sin 3t}{},$$

$$\underline{z(t)} = \mathcal{L}^{-1}[Z(p)] = \frac{-2e^{3t} \cos 3t + 4e^{3t} \sin 3t}{},$$

(7) $u_{tt} - c^2 u_{xx} = 0, \quad 0 \leq x < \infty$

$u(x, 0) = 0, \quad u_t(x, 0) = 0$ - initial conditions

$u(0, t) = f(t)$ - boundary condition.

$$\mathcal{L}[u_{tt}] = p^2 U(x, p) - \underbrace{p u(x, 0) - u_t(x, 0)}_{= 0 \text{ in this problem}}$$

$$\mathcal{L}[u_{xx}] = \int_0^\infty e^{-pt} \frac{\partial^2 u}{\partial x^2} dt = \frac{\partial^2}{\partial x^2} \int_0^\infty e^{-pt} u(x, t) dt$$

$$= \frac{\partial^2 U(x, p)}{\partial x^2}$$

Applying the Laplace transform to the wave equation, we have:

$$p^2 U(x, p) - c^2 \frac{\partial^2 U}{\partial x^2} = 0,$$

or $\frac{\partial^2 U}{\partial x^2} - \frac{p^2}{c^2} U = 0.$

This is an ordinary differential equation in x for $U(x,p)$, and its general solution is:

$$U(x,p) = A(p) e^{\frac{p}{c}x} + B(p) e^{-\frac{p}{c}x}$$

↑ ↑
arbitrary functions of p .

Since our solution should not diverge for $x \rightarrow \infty$, we must set $A(p) = 0$, so that

$$U(x,p) = B(p) e^{-\frac{p}{c}x}.$$

Using the initial condition $u(0,t) = f(t)$, and taking its Laplace transform,

$$U(0,p) = F(p), \quad \text{where } F(p) = \mathcal{L}[f],$$

we see that $B(p) = F(p)$.

Hence: $\underline{U(x,p) = F(p)e^{-\frac{p}{c}x}}$.

$$u(x,t) = \mathcal{L}^{-1}[U(x,p)] = \mathcal{L}^{-1}[F(p)e^{-\frac{x}{c}p}]$$

By the 2nd shift theorem, we have:

$$u(x,t) = f\left(t - \frac{x}{c}\right)$$

This answer is valid for $t - \frac{x}{c} \geq 0$, i.e. $x \leq ct$. For $t - \frac{x}{c} < 0$, $f(t - \frac{x}{c}) = 0$.

So, finally:

$$u(x,t) = \begin{cases} f\left(t - \frac{x}{c}\right), & x \leq ct \\ 0, & x > ct \end{cases}$$

This solution shows that the disturbance of the string at $x=0$ propagates along the string with velocity c .